

The effect of thermal non-equilibrium on kinetic nucleation

Sven Kiefer

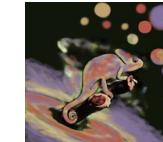
David Gobrecht, Leen Decin, Christiane Helling



ÖSTERREICHISCHE
AKADEMIE DER
WISSENSCHAFTEN



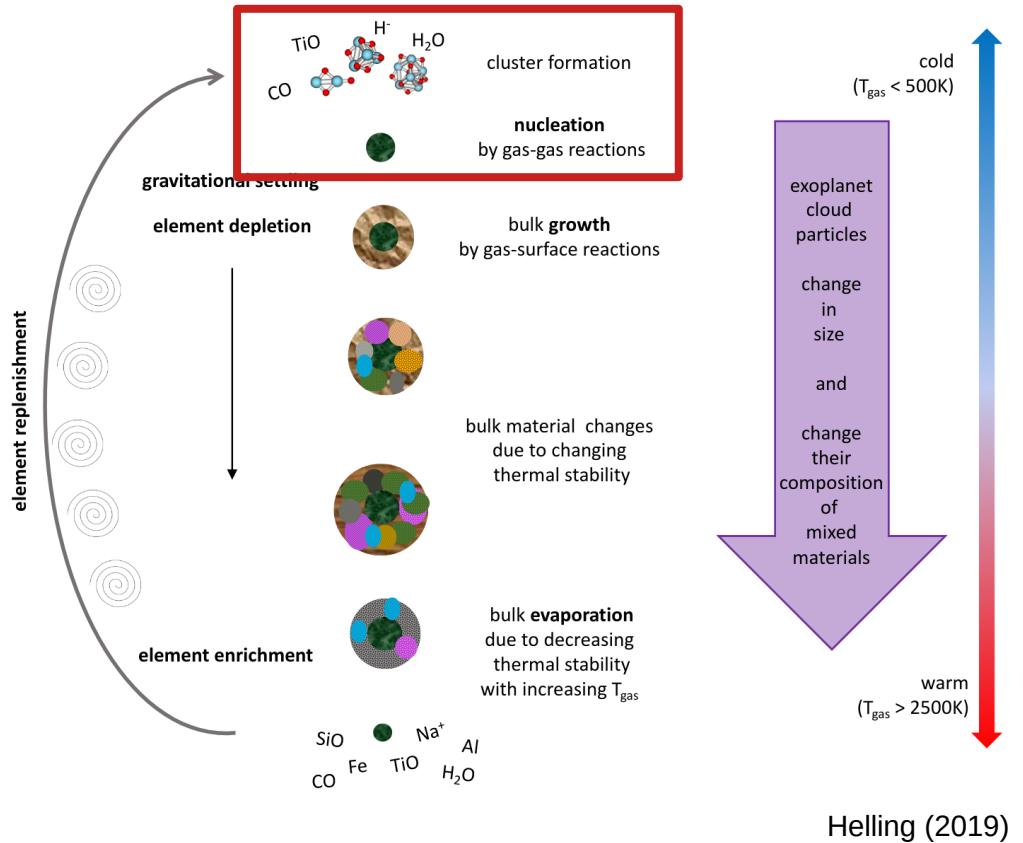
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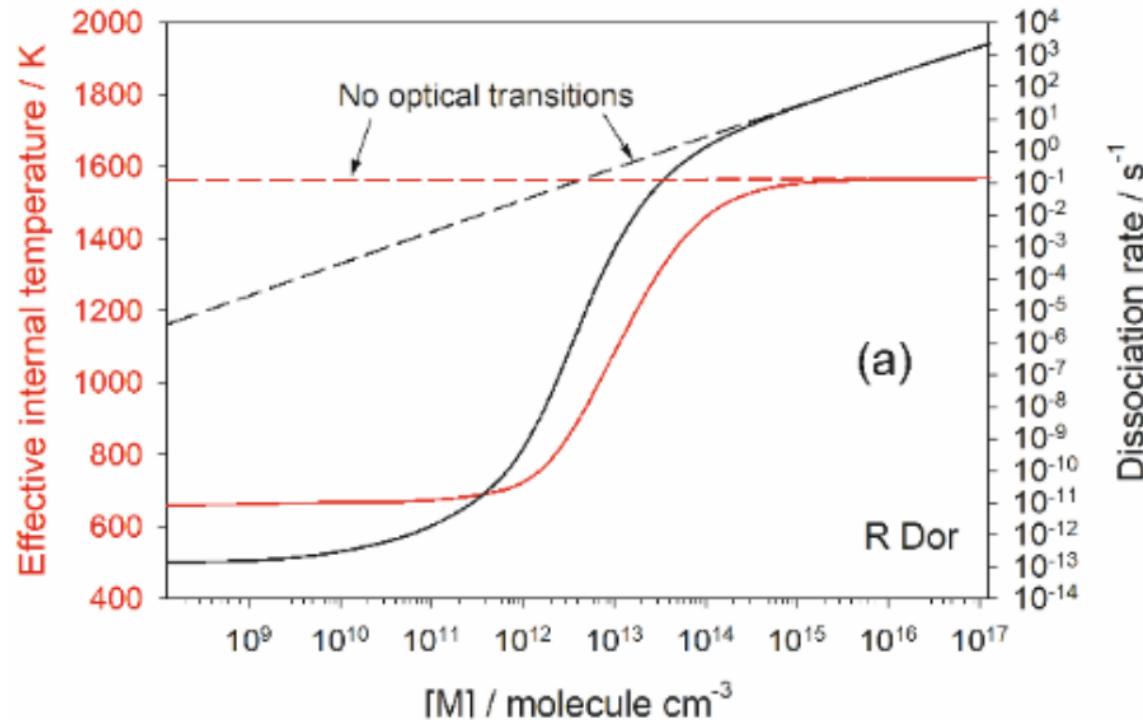
This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement no. 860470.

Overview

- Where to find thermal non-equilibrium
- How to nucleate
- Thermal non-equilibrium
- The effect of thermal non-equilibrium



Where to find thermal non-equilibrium



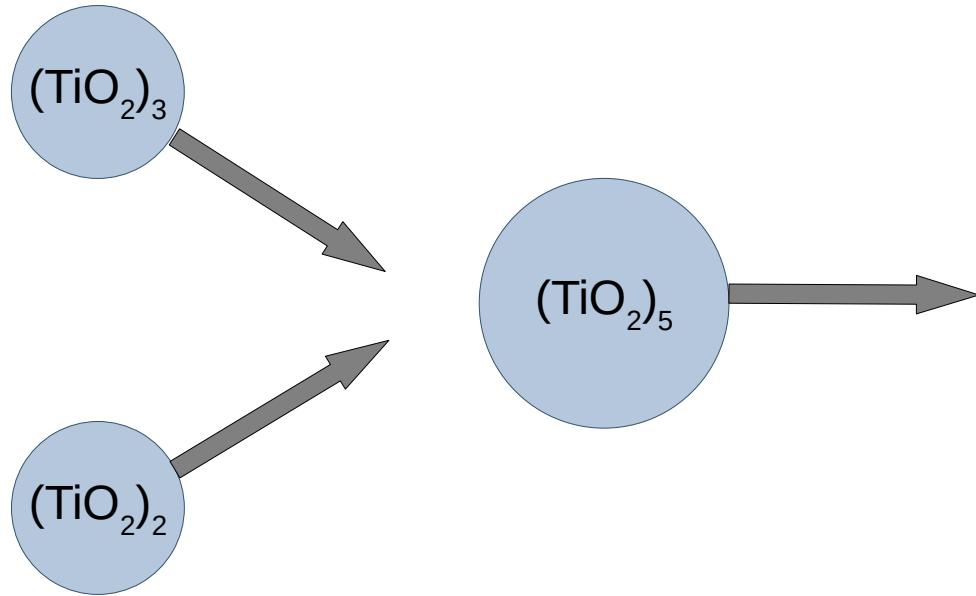
Plane and Robertson 2022:

- Outflows of AGB stars
- Dissociation of OSi(OH)₂
- Internal cooling via optical lines

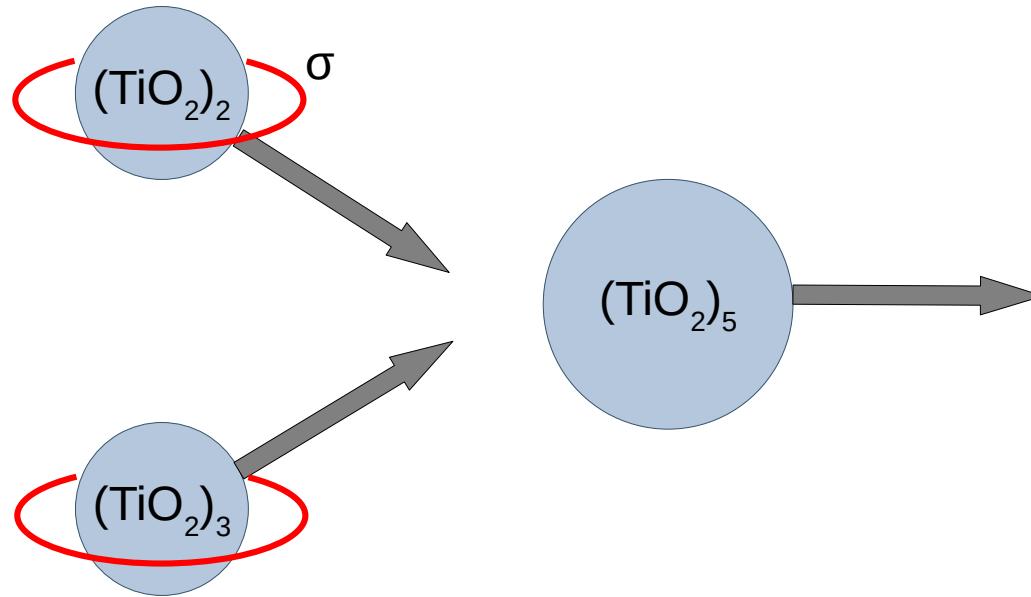
Observational evidence:

- Fonfria et al. 2008, 2017, 2021

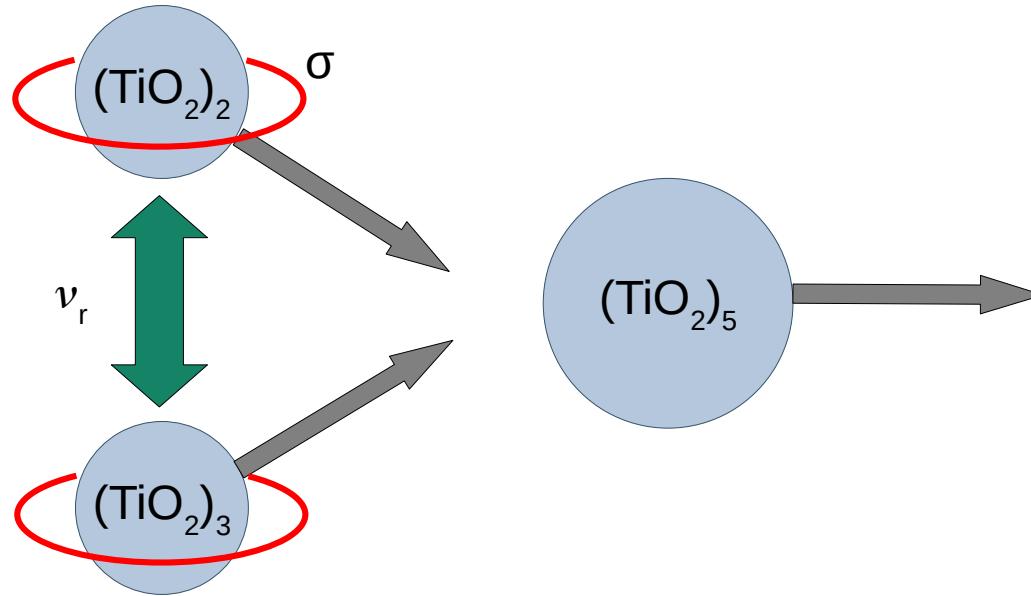
How to nucleate



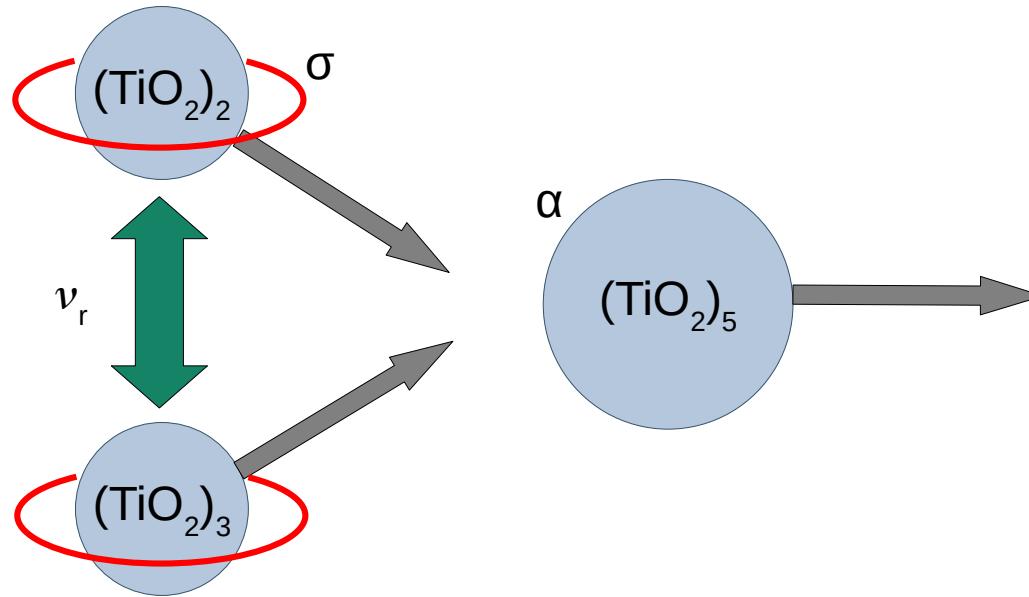
How to nucleate



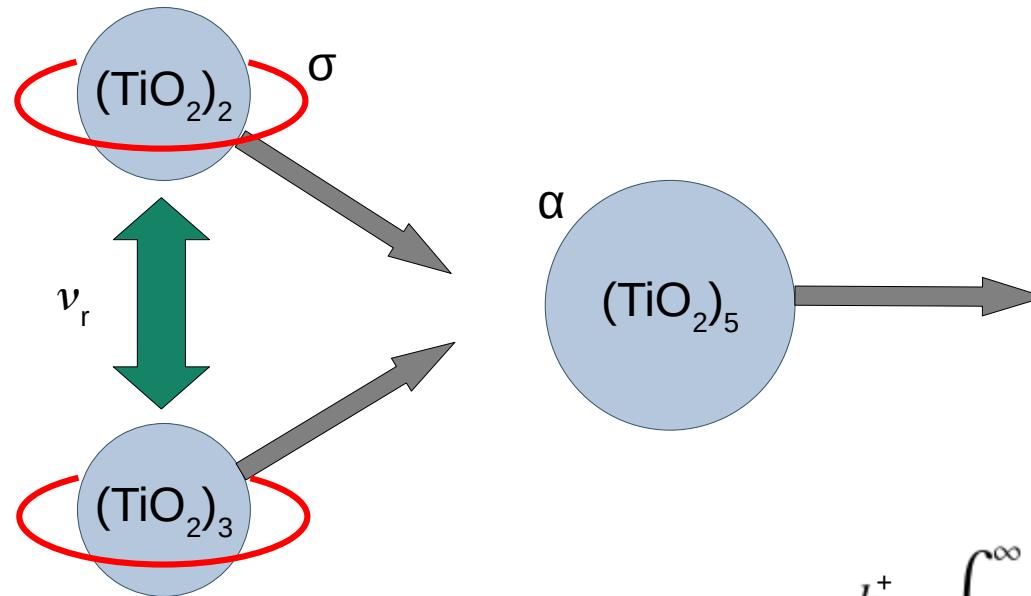
How to nucleate



How to nucleate

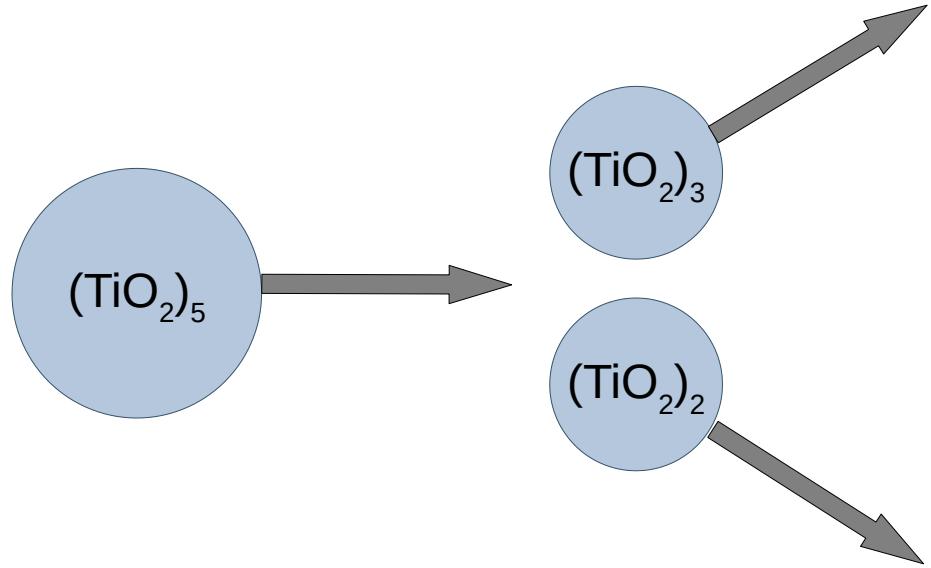


How to nucleate



$$k^+ = \int_0^\infty \alpha(v_r) \sigma(v_r) v_r f(v_r) dv_r$$

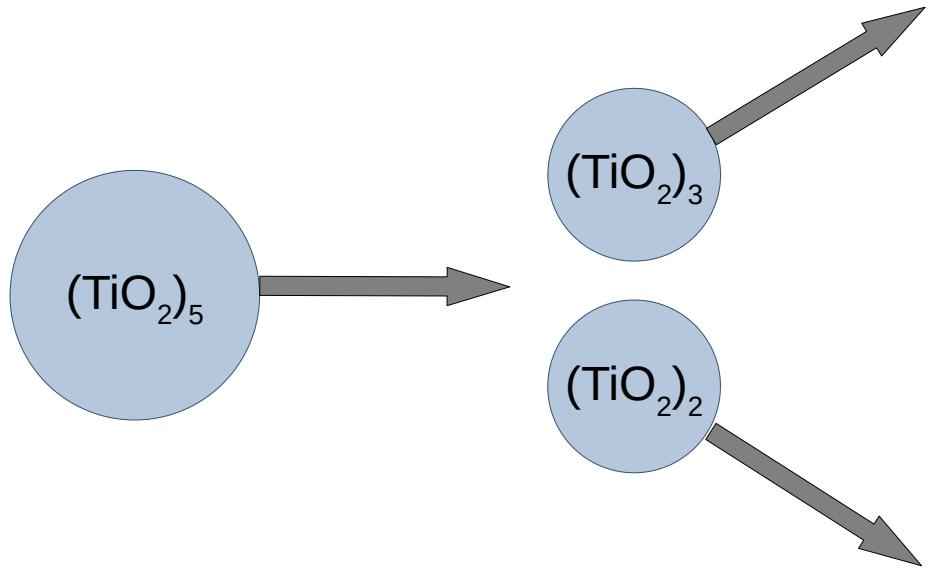
How to nucleate



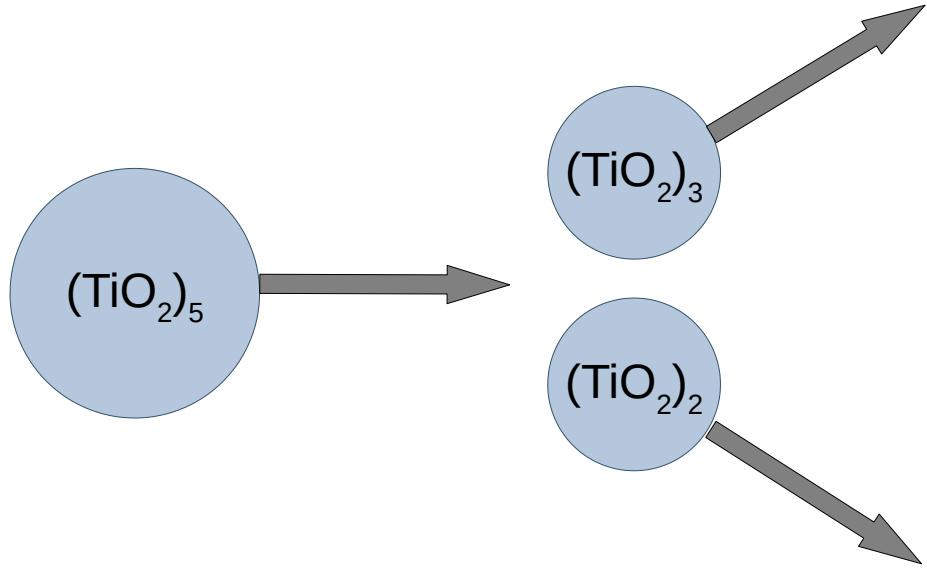
How to nucleate

Detailed balance in chemical equilibrium:

$$\frac{k_{N,M}^-}{k_{N,M}^+} = \frac{n_N^{\text{eq}} n_M^{\text{eq}}}{n_{N+M}^{\text{eq}}}$$



How to nucleate



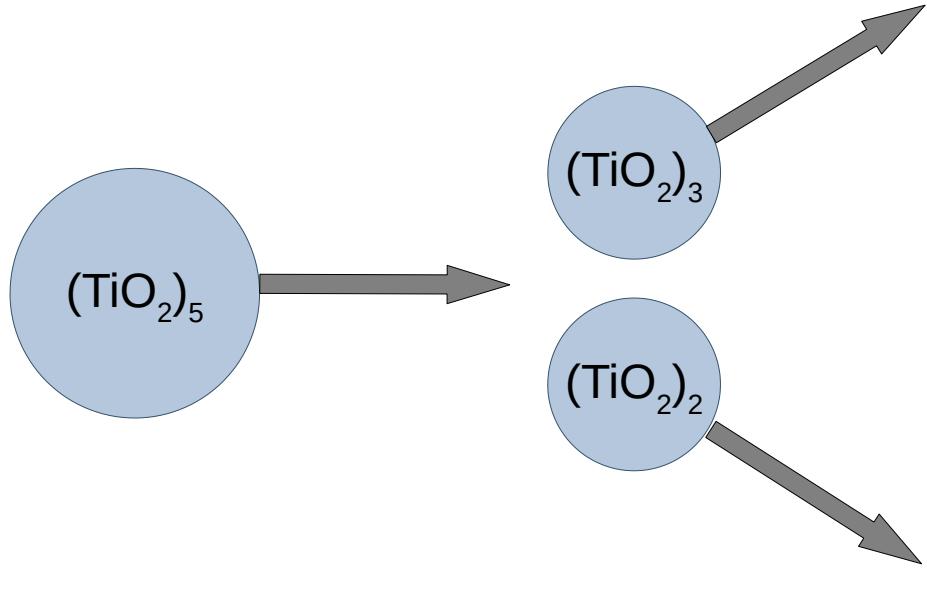
Detailed balance in chemical equilibrium:

$$\frac{k_{N,M}^-}{k_{N,M}^+} = \frac{n_N^{\text{eq}} n_M^{\text{eq}}}{n_{N+M}^{\text{eq}}}$$

Law of mass action in thermal equilibrium:

$$\frac{n_N^{\text{eq}} n_M^{\text{eq}}}{n_{N+M}^{\text{eq}}} = \frac{P^\circ}{k_B T} \exp\left(\frac{G_{N+M}^\circ - G_M^\circ - G_N^\circ}{k_B T}\right)$$

How to nucleate



Detailed balance in chemical equilibrium:

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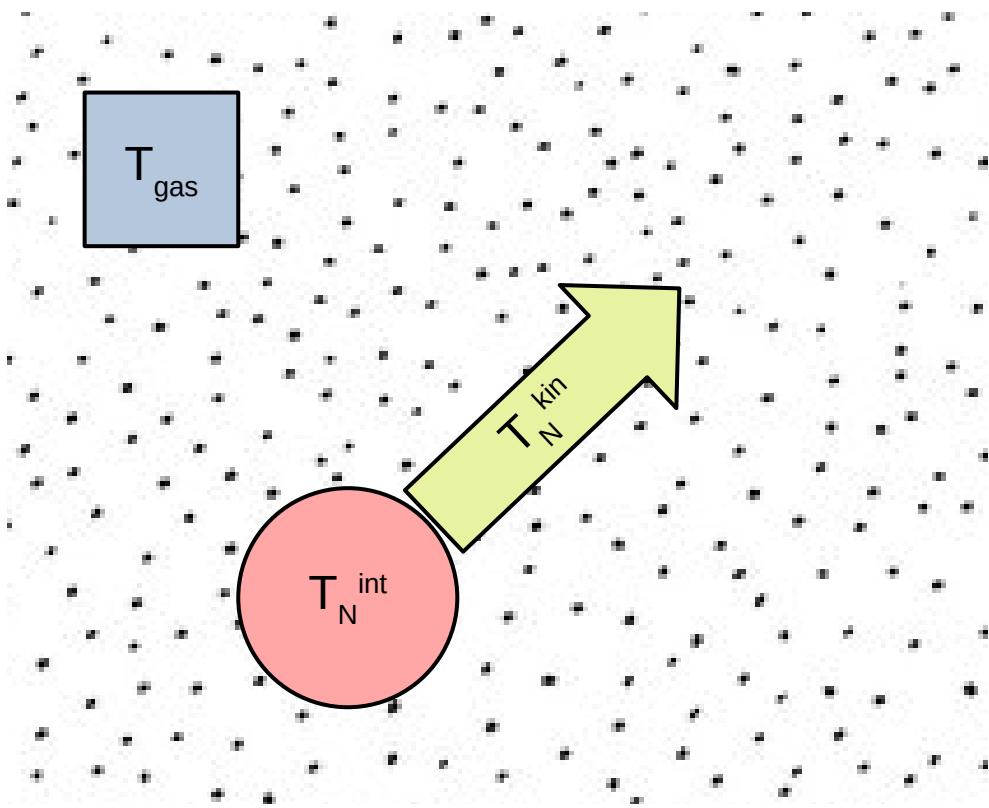
Law of mass action in thermal equilibrium:

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Backward rate:

$$k_{N,M}^- = k_{N,M}^+ \frac{P^\circ}{k_B T} \exp\left(\frac{G_{N+M}^\circ - G_M^\circ - G_N^\circ}{k_B T}\right)$$

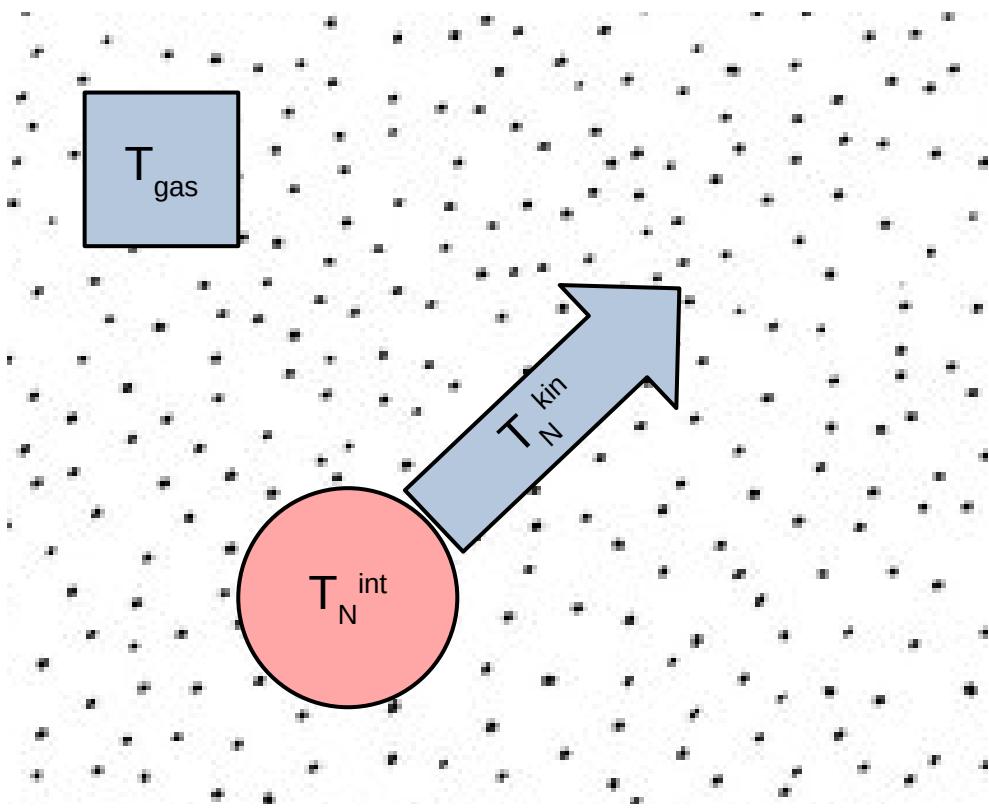
Thermal non-equilibrium



Types of thermal non-equilibrium:

- $T_{\text{gas}} \neq T_N^{\text{kin}} \neq T_N^{\text{int}}$
→ Kiefer et al. 2023

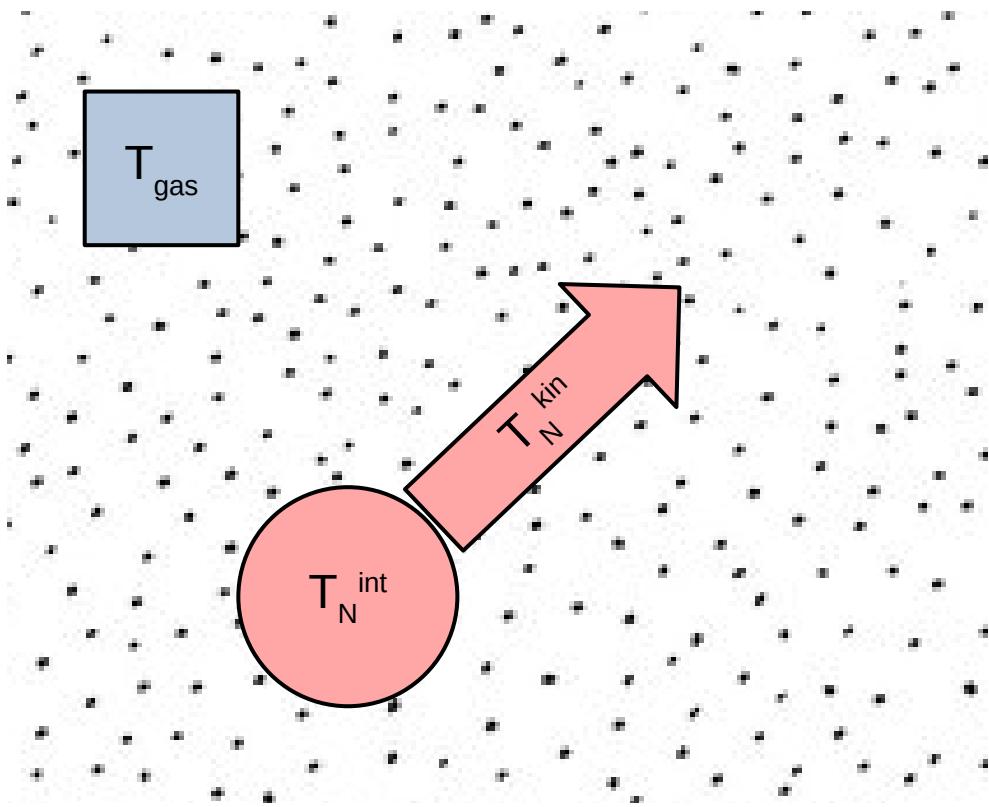
Thermal non-equilibrium



Types of thermal non-equilibrium:

- $T_{\text{gas}} \neq T_N^{\text{kin}} \neq T_N^{\text{int}}$
→ Kiefer et al. 2023
- $T_{\text{gas}} = T_N^{\text{kin}} \neq T_N^{\text{int}}$
→ Plane et al. 2022

Thermal non-equilibrium



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- $T_{\text{gas}} = T_N^{\text{kin}} \neq T_N^{\text{int}}$
→ Plane et al. 2022
- $T_{\text{gas}} \neq T_N^{\text{kin}} = T_N^{\text{int}}$
→ Patzer et al. 1998
→ Köhn et al. 2021

Thermal non-equilibrium

Thermal equilibrium

$$k^+ = \int_0^\infty \alpha(\nu_r) \sigma(\nu_r) \nu_r f(\nu_r) d\nu_r$$

$$k_{N,M}^- = k_{N,M}^+ \frac{P^\circ}{k_B T} \exp\left(\frac{G_{N+M}^\circ - G_M^\circ - G_N^\circ}{k_B T}\right)$$

Thermal non-equilibrium



Full derivation can be found in:
"The effect of thermal non-equilibrium on
kinetic nucleation" - Kiefer et al. 2023

1

Thermal non-equilibrium

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$$k^- = k^+ \frac{p^\circ}{kT_{\text{gas}}} A B C$$

$$A = \exp\left(\sum_{i \in \zeta} \frac{\delta(i)}{kT_i^{\text{kin}}} \left[G_i^\circ(T_{\text{gas}}) - iG_1^\circ(T_i^{\text{kin}}) + k(T_i^{\text{kin}} - T_{\text{gas}}) \right]\right)$$

$$B = \exp\left(\sum_{i \in \zeta} \frac{\delta(i)}{kT_i^{\text{kin}}} \omega_i(T_i^{\text{kin}}, T_i^{\text{int}})\right)$$

$$C = \left(\frac{kT_{\text{gas}} n_1^{\text{eq}}}{p^\circ}\right)^{-\sum_{i \in \zeta} \delta(i) i \frac{T_{\text{gas}}}{T_i^{\text{kin}}}}$$



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$$T^{\text{kin}} \neq T^{\text{int}}$$

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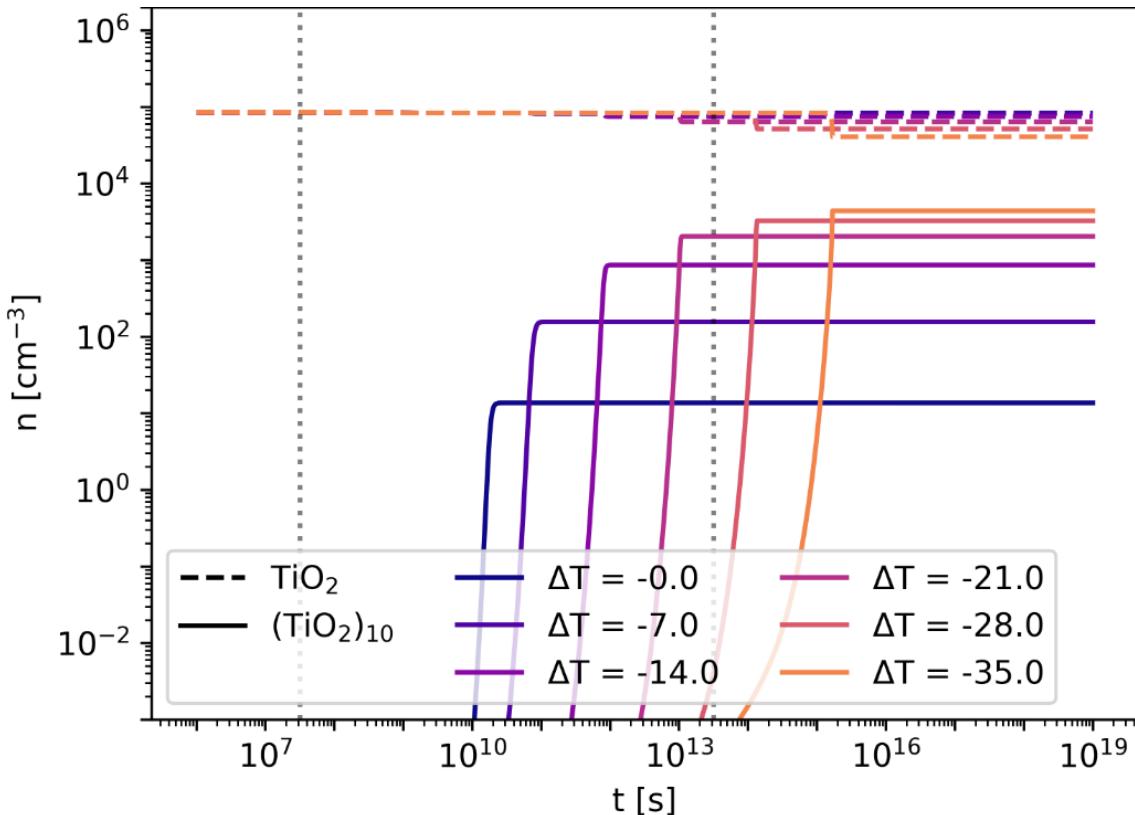
$$C = \left(\frac{kT_{\text{gas}} n_1^{\text{eq}}}{P^\circ} \right)^{-\sum_{i \in \zeta} \delta(i) i \frac{T_{\text{gas}}}{T_i^{\text{kin}}}}$$

$$T_{\text{gas}}^{\text{gas}} \neq T^{\text{kin}}$$

$$T^{\text{kin}} \neq T^{\text{int}}$$

Gauge

The effect of thermal non-equilibrium

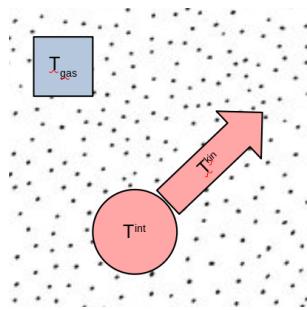


$$T_{\text{gas}} = 1250 \text{ K}$$

$$\Delta T = T_{(\text{TiO}_2)_{10}}^{\text{kin}} - T_{\text{TiO}_2}^{\text{kin}}$$

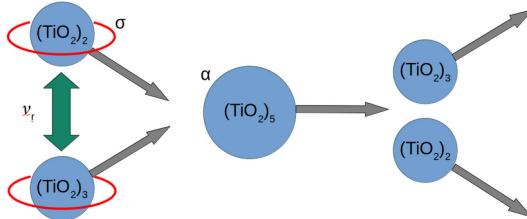
$$T_N^{\text{kin}} = T_{\text{gas}} + \frac{e^{N-1} - 1}{e^9 - 1} \Delta T$$

- Small temperature offsets can cause significant change in the number density of larger clusters.
- The change in number density depends on the thermal non-equilibrium present.
- Thermal non-equilibrium can both increase ($\Delta T < 0$) or decrease ($\Delta T > 0$) the formation of larger clusters.

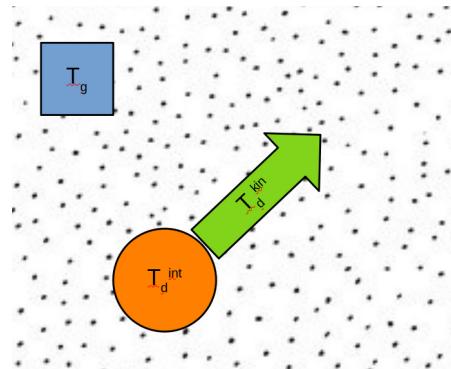


Summary

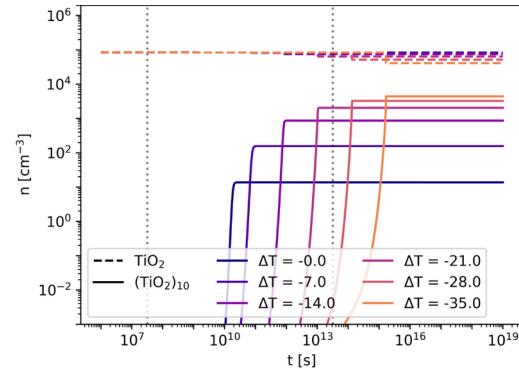
Kinetic nucleation can describe the formation of clusters under non-equilibrium conditions



Multiple types of thermal non-equilibrium are studied



Small temperature offsets can change the number density of clusters significantly

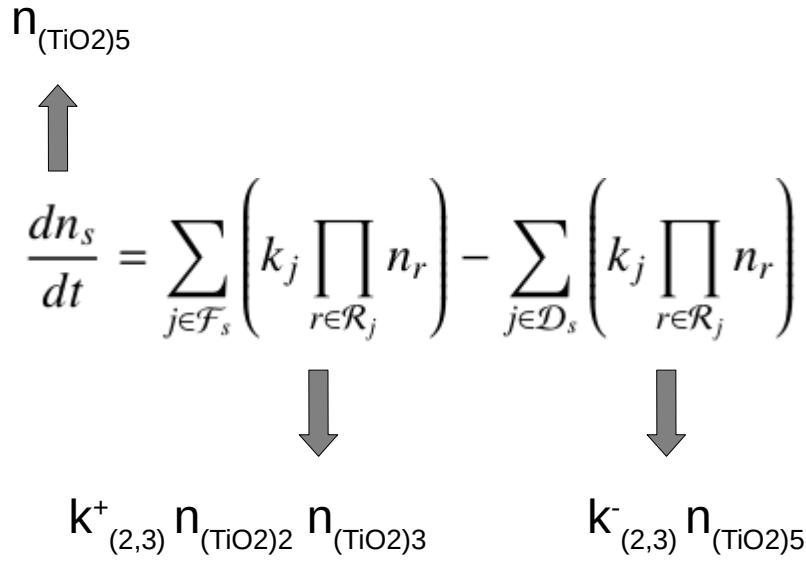


Check out the paper:
"The effect of thermal non-equilibrium on kinetic nucleation" - Kiefer et al. 2023



Additional Slides

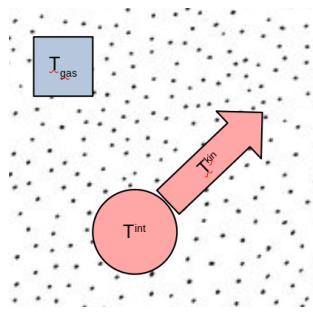
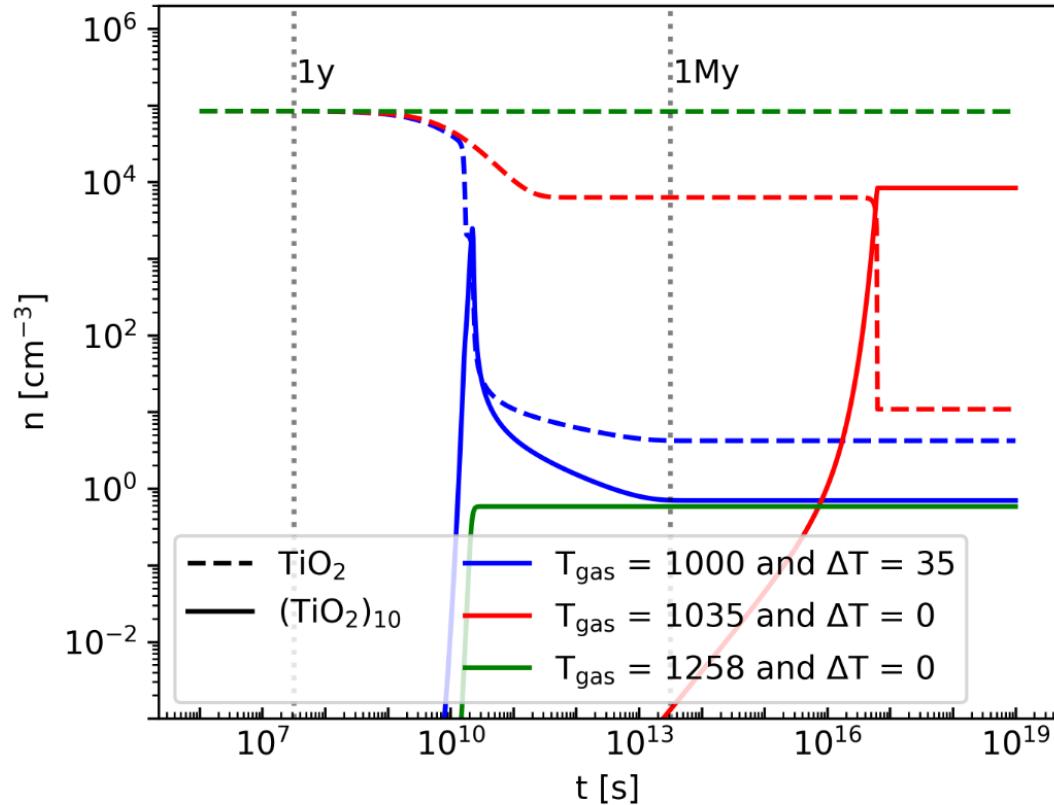
How to nucleate



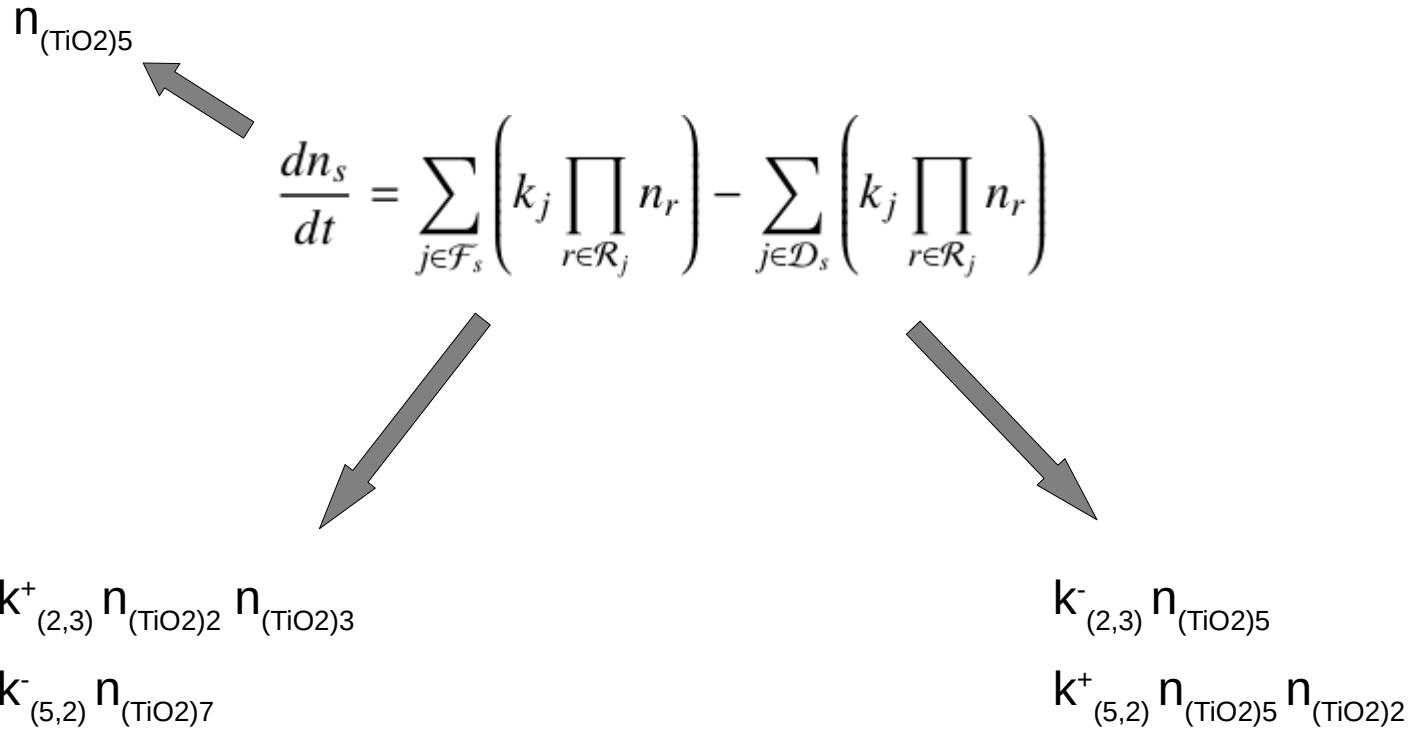
Nucleation network:

- Set of reactions that includes a formation path for larger cluster
- For reactions only involving clusters of size smaller than 4, we consider termolecular associations (3-Body reactions)
- For larger clusters, we consider bimolecular radiative associations (2-Body reactions)
- Backward rates are derived using detailed balance and the law of mass action

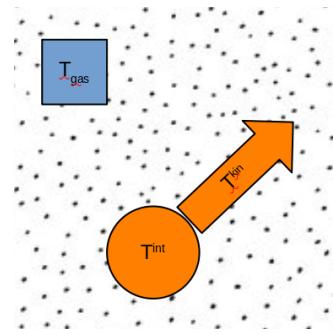
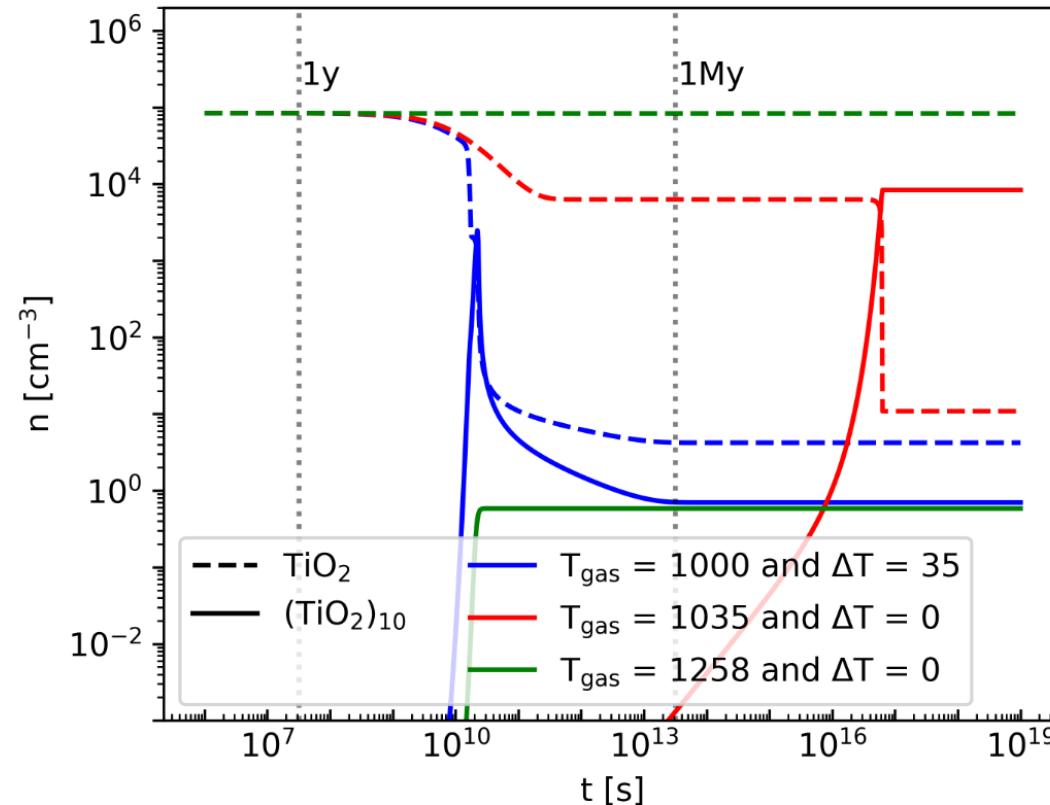
The effect of thermal non-equilibrium



How to nucleate



The effect of thermal non-equilibrium



Thermal adjustments

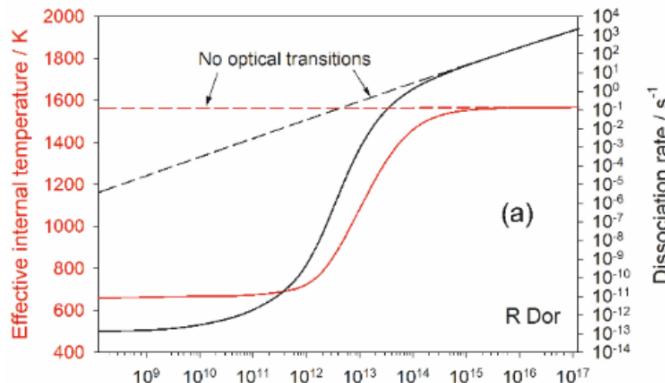
- Kinetic – Collisional

$$(\tau_{\text{gc}}^{\text{int}})^{-1} \approx \frac{2\bar{\alpha}_T}{D_f} n_{\text{gas}} r_N^2 \sqrt{\frac{8\pi k T_{\text{gas}}}{m_{\text{gas}}}}$$

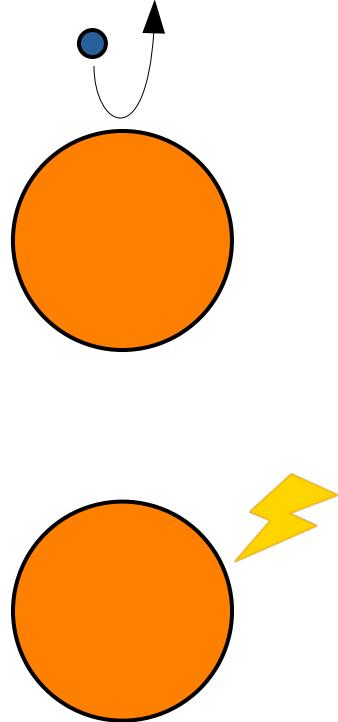
- Internal – Collisional

$$(\tau_{\text{gc}}^{\text{kin}})^{-1} \approx \frac{8m_{\text{gas}}}{3m_N} n_{\text{gas}} \pi r_N^2 \sqrt{\frac{8k T_{\text{gas}}}{\pi m_{\text{gas}}}}$$

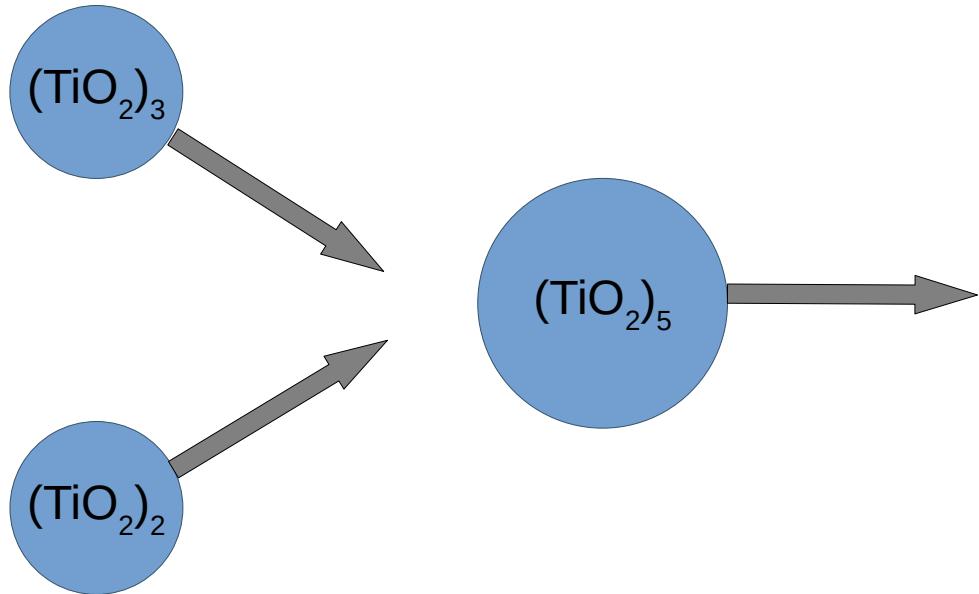
- Internal – Radiative



Plane et al. 2022



Inelastic collisions

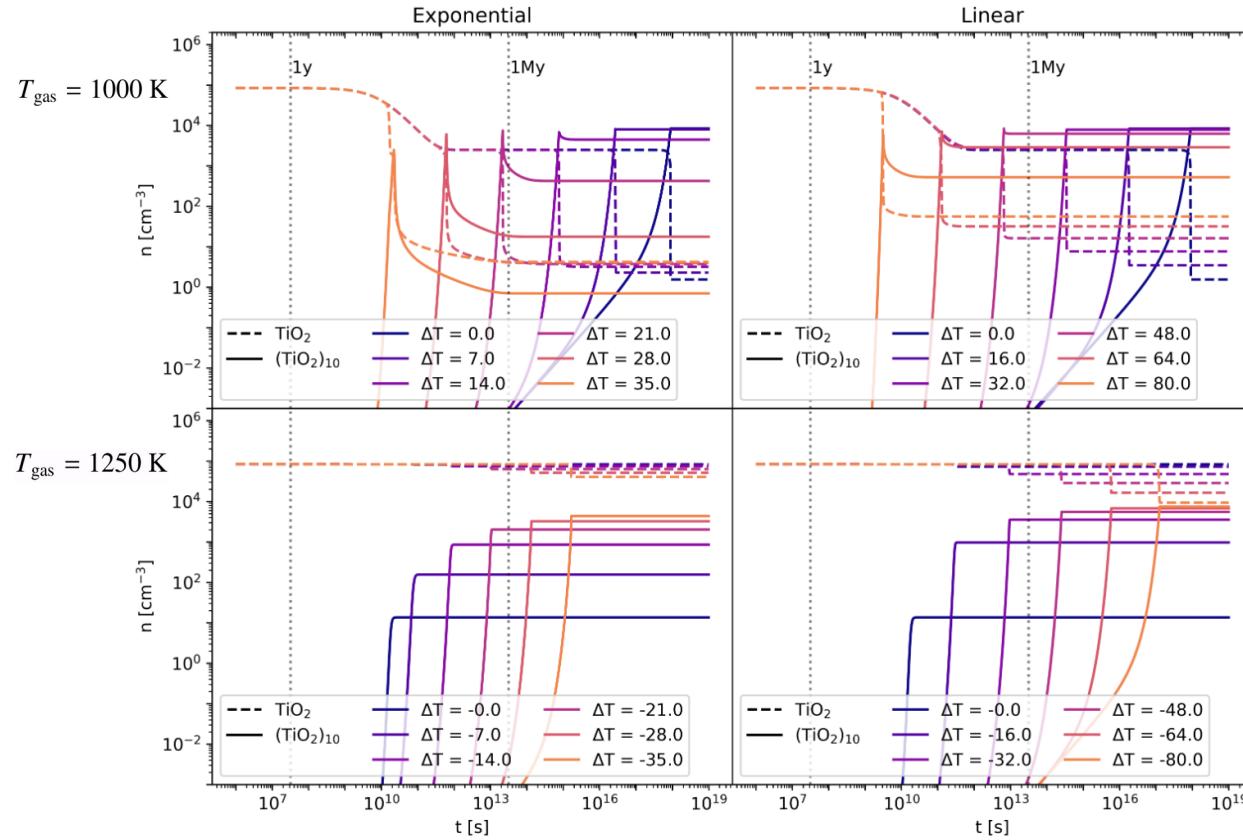


$$T^{\text{kin}} = \frac{(m_1 + m_2)}{\tilde{\mu}} = \frac{m_1 T_1^{\text{kin}} + m_2 T_2^{\text{kin}}}{(m_1 + m_2)} = T_1^{\text{kin}} + T_2^{\text{kin}} - \frac{\mu}{\mu_T}$$

$$\Delta E^{\text{kin}} = \frac{3k}{2}(T_N^{\text{kin}} + T_M^{\text{kin}} - T_{M+N}^{\text{kin}}) = \frac{3k}{2} \frac{\mu}{\mu_T}$$

$$\Delta T_{\text{inel}}^{\text{int}} = \frac{2}{D_f k} (\Delta E_{\text{inel}}^{\text{kin}} + \Delta E_{\text{frebind}}) = \frac{2}{D_f k} \left(\frac{3k\mu}{2\mu_T} + \Delta E_{\text{bind}} \right)$$

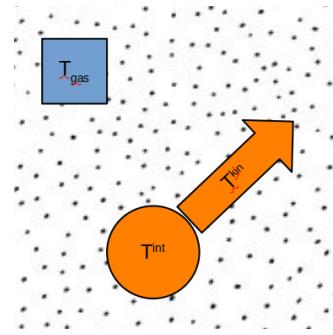
The effect of thermal non-equilibrium



$$\Delta T = T_{(\text{TiO}_2)_{10}}^{\text{kin}} - T_{\text{TiO}_2}^{\text{kin}}$$

$$T_N^{\text{kin}} = T_{\text{gas}} + \frac{e^{N-1} - 1}{e^9 - 1} \Delta T$$

$$T_N^{\text{kin}} = T_{\text{gas}} + \frac{(N-1)}{9} \Delta T$$



- Small temperature offsets can cause significant change in the number density of larger clusters.
- The change in number density depends on the thermal non-equilibrium present.
- Thermal non-equilibrium can both increase ($\Delta T < 0$) or decrease ($\Delta T > 0$) the formation of larger clusters.

Derivation of the Backward rate

$$G^{non-eq}(T_0^{\text{int}}, \dots, T_r^{\text{int}}, T_0^{\text{kin}}, \dots, T_r^{\text{kin}}, p_0, \dots, p_r, N_0, \dots, N_r) = \sum_{i=0}^r G_i^{non-eq}(T_i^{\text{int}}, T_i^{\text{kin}}, p_i, N_i)$$

$$= \sum_{i=0}^r G_i(T_i^{\text{kin}}, p_i, N_i) + N_i \omega_i(T_i^{\text{kin}}, T_i^{\text{int}}),$$

$$C = \sum_{i=1}^r i N_i$$

$$G = N\mu$$

$$\begin{aligned} \mathcal{L} &= N_{\text{gas}}\mu_{\text{gas}}(T_{\text{gas}}, p) + N_{\text{gas}}kT_{\text{gas}} \ln\left(\frac{N_{\text{gas}}}{N}\right) - \lambda C \\ &+ \sum_{i=1}^r N_i\mu_i(T_i^{\text{kin}}, p) + N_i kT_i^{\text{kin}} \ln\left(\frac{N_i}{N}\right) \\ &+ N_i \omega_i(T_i^{\text{kin}}, T_i^{\text{int}}) + \lambda i N_i. \end{aligned}$$

$$k^- = \frac{k^+ p^\ominus}{kT_{\text{gas}}} \exp\left(\sum_{i \in \zeta} \frac{\delta(i)}{kT_i^{\text{kin}}} [\mu_i^\ominus(T_i^{\text{kin}}) - i\mu_1^\ominus(T_{\text{gas}}) + k(T_i^{\text{kin}} - T_{\text{gas}}) + \omega_i(T_i^{\text{kin}}, T_i^{\text{int}})]\right) \left(\frac{kT_{\text{gas}} \dot{n}_1}{p^\ominus}\right)^{-\sum_{i \in \zeta} \delta(i)i T_{\text{gas}}/T_i^{\text{kin}}}$$

$$k_j^- = k_j^+ \frac{\dot{n}_A \dot{n}_B}{\dot{n}_{A+B}}$$

$$\frac{\dot{N}_j}{\dot{N}_{\text{gas}}} = \exp\left(\frac{-\mu_j(T_j^{\text{kin}}, p)}{kT_j^{\text{kin}}}\right) \exp\left(\frac{-k(T_j^{\text{kin}} - T_{\text{gas}}) - \omega_j(T_j^{\text{kin}}, T_j^{\text{int}}) - \lambda j}{kT_j^{\text{kin}}}\right).$$

$$\lambda = -\mu_1(T_{\text{gas}}, p) - kT_{\text{gas}} \ln\left(\frac{N_1}{N}\right)$$

$$N_{\text{gas}} \gg \sum_{i=1}^r N_i$$

$$N_{\text{gas}} T_{\text{gas}} \gg \sum_{i=1}^r N_i T_i^{\text{kin}}$$

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial N_j} &= \mu_j(T_j^{\text{kin}}, p) + kT_j^{\text{kin}} \ln\left(\frac{N_j}{N}\right) + \omega_j(T_j^{\text{kin}}, T_j^{\text{int}}) \\ &+ \lambda j x_1 + kT_j^{\text{kin}} \frac{N - N_j}{N} - \sum_{\substack{i=0 \\ i \neq j}}^r \frac{N_i kT_i^{\text{kin}}}{N}, \\ &\approx \mu_j(T_j^{\text{kin}}, p) + kT_j^{\text{kin}} \ln\left(\frac{N_j}{N}\right) + \omega_j(T_j^{\text{kin}}, T_j^{\text{int}}) \\ &+ \lambda j + k(T_j^{\text{kin}} - T_{\text{gas}}), \end{aligned}$$

Maxwell-Boltzmann

$$M_T \equiv \frac{m_1}{T_1^{\text{kin}}} + \frac{m_2}{T_2^{\text{kin}}} = \frac{m_1 T_2^{\text{kin}} + m_2 T_1^{\text{kin}}}{T_1^{\text{kin}} T_2^{\text{kin}}}$$

$$\mu \equiv \frac{m_1 m_2}{m_1 + m_2}$$

$$\mu_T \equiv \frac{\frac{m_1}{T_1^{\text{kin}}} \frac{m_2}{T_2^{\text{kin}}}}{\frac{m_1}{T_1^{\text{kin}}} + \frac{m_2}{T_2^{\text{kin}}}} = \frac{m_1 m_2}{m_1 T_2^{\text{kin}} + m_2 T_1^{\text{kin}}}$$

$$\boldsymbol{\nu}_r \equiv \boldsymbol{\nu}_1 - \boldsymbol{\nu}_2$$

$$\boldsymbol{\nu}_T \equiv \frac{\frac{m_1}{T_1^{\text{kin}}} \boldsymbol{\nu}_1 + \frac{m_2}{T_2^{\text{kin}}} \boldsymbol{\nu}_2}{\frac{m_1}{T_1^{\text{kin}}} + \frac{m_2}{T_2^{\text{kin}}}} = \frac{m_1 T_2^{\text{kin}} \boldsymbol{\nu}_1 + m_2 T_1^{\text{kin}} \boldsymbol{\nu}_2}{m_1 T_2^{\text{kin}} + m_2 T_1^{\text{kin}}}$$



$$\begin{aligned} \int_{\mathbb{R}^3} f_r(\boldsymbol{\nu}_r) d\boldsymbol{\nu}_r &= \int_{\mathbb{R}^3} f(\boldsymbol{\nu}_1) d\boldsymbol{\nu}_1 \int_{\mathbb{R}^3} f(\boldsymbol{\nu}_2) d\boldsymbol{\nu}_2 \\ &= \left(\frac{1}{2\pi k} \right)^3 \left(\frac{m_1 m_2}{T_1^{\text{kin}} T_2^{\text{kin}}} \right)^{3/2} \\ &\quad \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \exp \left(-\frac{M_T \boldsymbol{\nu}_T^2 + \mu_T \boldsymbol{\nu}_r^2}{2k} \right) d\boldsymbol{\nu}_r d\boldsymbol{\nu}_T \\ &= \int_{\mathbb{R}^3} \left(\frac{\mu_T}{2\pi k} \right)^{3/2} \exp \left(-\frac{\mu_T \boldsymbol{\nu}_r^2}{2k} \right) d\boldsymbol{\nu}_r. \end{aligned}$$



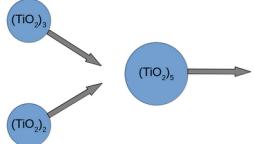
$$f_r(\boldsymbol{\nu}_r) d\boldsymbol{\nu}_r = \left(\frac{\mu_T}{2\pi k} \right)^{3/2} 4\pi \boldsymbol{\nu}_r^2 \exp \left(-\frac{\mu_T \boldsymbol{\nu}_r^2}{2k} \right) d\boldsymbol{\nu}_r$$



$$\frac{m_1}{T_1^{\text{kin}}} \boldsymbol{\nu}_1^2 + \frac{m_2}{T_2^{\text{kin}}} \boldsymbol{\nu}_2^2 = M_T \boldsymbol{\nu}_T^2 + \mu_T \boldsymbol{\nu}_r^2$$

$$k_j^+ = \int_0^\infty \pi (r_1 + r_2)^2 \nu_r \left(\frac{\mu_T}{2\pi k} \right)^{3/2} 4\pi \nu_r^2 \exp \left(-\frac{\mu_T \nu_r^2}{2k} \right) d\nu_r = \pi (r_1 + r_2)^2 \sqrt{\frac{8k}{\pi \mu_T}}$$

- Temperature-weighted reduced mass μ_T
- In thermal equilibrium $\mu_T = \mu / T$



Inelastic collisions

$$\begin{aligned}\tilde{\mathbf{v}} &\equiv \frac{m_1}{m_1 + m_2} \mathbf{v}_1 + \frac{m_2}{m_1 + m_2} \mathbf{v}_2 \\ \tilde{\mathbf{v}}_T &\equiv \frac{\frac{m_2}{T_1^{\text{kin}}} m_1 \mathbf{v}_1 - \frac{m_1}{T_2^{\text{kin}}} m_2 \mathbf{v}_2}{\frac{m_2}{T_1^{\text{kin}}} + \frac{m_1}{T_2^{\text{kin}}}} \\ \tilde{\mu} &\equiv \frac{(m_1 + m_2)^2}{m_1 T_1^{\text{kin}} + m_2 T_2^{\text{kin}}} \\ \tilde{M} &\equiv \frac{\frac{m_2}{T_1^{\text{kin}}} + \frac{m_1}{T_2^{\text{kin}}}}{m_1 m_2}\end{aligned}$$

$$\frac{m_1}{T_1^{\text{kin}}} \mathbf{v}_1^2 + \frac{m_2}{T_2^{\text{kin}}} \mathbf{v}_2^2 = \tilde{M} \tilde{\mathbf{v}}_T^2 + \tilde{\mu} \tilde{\mathbf{v}}^2$$

$$d\mathbf{v}_1 d\mathbf{v}_2 = \frac{1}{\mu^3} d\tilde{\mathbf{v}}_T d\tilde{\mathbf{v}}.$$

$$\begin{aligned}\iint_{\mathbb{R}^3} f(\tilde{\mathbf{v}}) d\tilde{\mathbf{v}} &= \iint_{\mathbb{R}^3} f(\mathbf{v}_1) f(\mathbf{v}_2) d\mathbf{v}_1 d\mathbf{v}_2 \\ &= \iint_{\mathbb{R}^3} \left(\frac{1}{2\pi k} \right)^3 \left(\frac{m_1 m_2}{T_1 T_2} \right)^{3/2} \exp\left(-\frac{\tilde{M} \tilde{\mathbf{v}}_T^2 + \tilde{\mu} \tilde{\mathbf{v}}^2}{2k}\right) \frac{1}{\mu^3} d\tilde{\mathbf{v}}_T d\tilde{\mathbf{v}} \\ &= \int_{\mathbb{R}^3} \left(\frac{\tilde{\mu}}{2\pi k} \right)^{3/2} \exp\left(-\frac{\tilde{\mu} \tilde{\mathbf{v}}^2}{2k}\right) d\tilde{\mathbf{v}}\end{aligned}$$

$$T^{\text{kin}} = \frac{(m_1 + m_2)}{\tilde{\mu}} = \frac{m_1 T_1^{\text{kin}} + m_2 T_2^{\text{kin}}}{(m_1 + m_2)} = T_1^{\text{kin}} + T_2^{\text{kin}} - \frac{\mu}{\mu_T}$$

$$(m_1 + m_2) \tilde{\mathbf{v}} = m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2$$

$$\Delta T_{\text{inel}}^{\text{int}} = \frac{2}{D_f k} (\Delta E_{\text{inel}}^{\text{kin}} + \Delta E_{\text{frebind}}) = \frac{2}{D_f k} \left(\frac{3k\mu}{2\mu_T} + \Delta E_{\text{bind}} \right)$$

Time dependency

$$T_i(\Delta t) = T_{\text{gas}} + T_{i,0} \exp(-\Delta t/\tau_i)$$



$$k_j^+(\Delta t_1, \Delta t_2) = \pi(r_1 + r_2)^2 \sqrt{\frac{8k}{\pi m_1 m_2} (m_2 T_1^{\text{kin}}(\Delta t_1) + m_1 T_2^{\text{kin}}(\Delta t_2))}$$



$$P_i(\Delta t) = (\tau_i)^{-1} \exp(-\Delta t/\tau_i)$$



$$\tilde{k}_j^+ = \int_0^\infty \int_0^\infty P_1(\Delta t_1) P_2(\Delta t_2) k_j^+(\Delta t_1, \Delta t_2) d(\Delta t_1) d(\Delta t_2)$$

$$\tau_{A,j}^{\text{coll},+} = \frac{1}{n_B k_j^+}$$



$$\tau_{A,j}^{\text{coll},+} = \tau_{\text{gc}}^{\text{kin}}$$



$$\epsilon_B = \frac{8}{3} \frac{r_B^2}{(r_A + r_B)^2} \sqrt{\frac{m_{\text{gas}} m_B T_{\text{gas}}}{m_A (m_A T_B + m_B T_A)}}$$



$$\epsilon_{(\text{TiO}_2)_2} = 0.053 \sqrt{\frac{T_{\text{gas}}}{T_{(\text{TiO}_2)_2}^{\text{kin}}}}$$